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CHARGE-TRANSFER INTERACTIONS BETWEEN

TRANSITION METAL HEXAFLUORIDES AND XENON.

by

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Charge transfer transitions are reported for MF ₆ (M = W, Mo, U, Re, Ir) with			
Xe. The nature of the charge transfer complexes is discussed. Based on			
spectroscopic data new estimates of ${ m MF}_6$ electron a	affinities are presented		
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I. INTRODUCTION

In the course of spectroscopic studies of paramagnetic transition metal hexafluorides, it became apparent that additional information might be obtained if a high symmetry host material for mixed crystals could be found. Such systems would be particularly useful for the study of the r_{8g} (0_h^{\star}) states, in which the Jahn-Teller interaction is of particular interest. In the assumption of the Yenon might serve as a good host; it has no crystal vibrational frequencies in the range of intramplecular MF₆ vibrations, and size and orientation considerations indicate that MF₆ might go into the Xe lattice substitutionally. However, it was found that when IrF_6 , which is yellow, is dissolved in liquid Xe, the solution is totally opaque, although purple in reflection. It is the purpose of this paper to demonstrate that the new absorption band for the IrF_6 -Xe system is due to a low-lying intermolecular charge transfer (CT) transition between Xe and MF₆ (Xe + MF₆ $\xrightarrow{h\nu}$ Xe $^+$ MF₅.) Such CT transitions and associated complexes with MF₆ are known, but have not been previously observed for the rare gases.

Several experiments have been carried out to verify these conclusions and to increase understanding of this phenomenon. The absorption spectra of MF_6/Xe (M = Ir, Re, W, Mo, U) at liquid nitrogen temperature have been taken. An IrF_6/Kr sample and a gas phase IrF_6/Xe sample were also prepared and investigated. Besides demonstrating the CT nature of the new transitions, these experiments also give valuable information on the exceptionally high electron affinities of the hexafluorides. Observation of "local" (intramolecular) IrF_6 transitions in the near IR also permits conclusions to be drawn concerning stability of the IrF_6-Xe complex.

II. EXPERIMENTAL

Handling of hexafluorides has been previously described. Research grade Xenon and Krypton (Linde) were used and were further purified by distillation to remove any traces of H_2O , a very serious impurity for the hexafluorides.

Crystals were grown from the melt by suspending the sample cell a few centimeters above the surface of liquid nitrogen in a closed dewar. Although crystals grown this rapidly (~20 min.) are certainly not high quality single crystals, they are of adequate quality to allow spectra to be taken. Visible and near UV absorption spectra were taken on a McPherson 285 monochromator with photoelectric detection. Near IR spectra were obtained on a McPherson 2051 with a 77K InAs (Texas Instruments) detector. Some preliminary spectra were also obtained on a Cary 17.

III. THEORY

The theory of CT transitions and complexes is well known and will be outlined only briefly here. For 1:n complexes (in this case n is either 12Xe or 1Xe), the energy of the CT transition is:

$$hv_{CT} = I_d - E_a + (G_1 - n'G_0) + (X_1 - n'X_0)$$
 (1)

in which

 I_d = Ionization potential of the donor D (Xe).

 $E_a = Electron affinity of the acceptor A (MF₆).$

 G_1 = "Normal" interaction of D^+ and A^- , specifically neglecting the CT interactions.

 G_0 = "Normal" interaction of D and A.

 X_1 = Additional interaction between D^+ and A^- due to proximity of D-A configuration.

 \mathbf{X}_0 = Additional interaction between D and A due to proximity of the $\mathbf{D}^+\mathbf{A}^-$ configuration.

n = Number of donors in the complex.

n' = In the limit of weak complexes, n' = n. For stronger complexes n' < n due to saturation effects.⁵

 $\mathbf{X}_{\mathbf{0}}$ can be approximated by second order perturbation theory, as

$$X_0 \sim -\frac{\beta_0^2}{\Lambda} , \qquad (2)$$

for which

$$\Delta = I_d - E_a + (G_1 - n'G_0)$$

 $\beta_0 = \langle \psi(D,A) | \mathcal{H} | \psi(D^+,A^-) \rangle.$

For the purpose of estimating electron affinities, the following approxi-

mate equation is useful:

$$hv_{CT} = I_d - E_a + G_1.$$
 (3)

Location and character of excited electronic states of the donor and acceptor are of importance in more detailed considerations. The excited electronic states of Xe are so high in energy as to have a negligible effect in this regard and will be ignored here. The onset of intramolecular CT transitions for the hexafluorides is given in Table 1. The lowest energy ligand-field states of ${\rm Ir} F_6$ are shown in Figure 1. Note that in the above considerations solvent effects have been neglected.

IV. RESULTS

The onset frequencies of the intermolecular CT transition for 1/2% MF_6/Xe ($\epsilon=1$ cm) are listed in Table 1. The observed transitions are found to be broad (>10,000 cm⁻¹) featureless bands, as expected of a CT transition. The frequencies of the near IR transitions of 0.1% IrF_6/Xe ($\epsilon=1$ cm) are given in Table 2; a comparison of IrF_6/Xe data and neat IrF_6 is given in Table 3. The intensities of the near IR transitions of IrF_6/Xe are estimated to be enhanced by two orders of magnitude over those of neat IrF_6 . The band widths are increased by an order of magnitude over the observed band widths in neat IrF_6 . Near IR and visible spectra of gas phase 1/2% IrF_6/Xe (10 torr-liter Xe, V = 10 ml, T = 300K, $\epsilon=1$ cm) and solid 0.01% IrF_6/Kr ($\epsilon=1$ cm) show only the intramolecular CT transition at 20,000 cm⁻¹ characteristic of IrF_6 . Freezing the yellow gas phase IrF_6/Xe mixture in liquid nitrogen gives an opaque-purple solid; warming the sample to the melting point of Xe gives an opaque-purple liquid which upon warming evaporates, leaving yellow solid IrF_6 . With the evaporation of IrF_6 , a pale yellow gas obtains, leaving no residue behind.

V. DISCUSSION

Plausibility of the intermolecular CT transition hypothesis (MF $_6$ + Xe hv [4F $_6$ Xe $^+$]) can be easily demonstrated. PtF $_6$ is known to chemically react with Xe, 6 the first step in the reaction being

$$PtF_6 + Xe \rightarrow PtF_6^-Xe^+$$
.

Since PtF₆ has a higher electron affinity (215 kcal/mole⁷) than the hexafluorides studied here, it is not unreasonable to suspect that the MF₆⁻Xe⁺ electronic configuration would be an excited state rather than the ground state as in PtF₆/Xe. A rough calculation can be used to demonstrate this point quite adequately. The frequency of the CT transition in IrF₆/Xe using Eq. (3) with I_d = 281 kcal/mole, E_a (IrF₆) \geq 136 kcal/mole , G₁ --70 kcal/mole, is found to be $h\nu_{CT} \leq 26,000 \text{ cm}^{-1}$.

Experimental verification of the CT transition hypothesis is given in Table 1. This data indicates that E_a (IrF_6) > E_a (MoF_6) > E_a (WF_6) which agrees with the expected trends. Estimates of the electron affinities of MF_6 are made in Table 1. The best procedure for estimating these values seems to be to use E_a (WF_6) as determined by the collisional ionization method and to estimate the others through the difference in the onset frequency of the intermolecular CT transition. The main source of uncertainty in this method is neglect of the interaction of the MF_6 -Xe electronic configuration with the ground (MF_6 Xe) and locally-excited (MF_6 Xe) configurations. The latter might be more significant for IrF_6 /Xe for which the CT transition occurs in the midst of the low-lying IrF_6 ligand field transitions (see Figure 1) and -10,000 cm⁻¹ from the intramolecular CT band of IrF_6 .

Since CT transitions often imply the formation of a CT complex, it is of

interest to determine if the MF $_6$ -Xe system forms such complexes. The CT transition in ${\rm IrF}_6$ /Xe is the closest to the ground state; thus, it is expected that the ${\rm IrF}_6$ -Xe complex would be the most stable. The usual approach for determining stability of such complexes employs the method of Benesi and Hildebrand to find the equilibrium constant for complex formation. However, the near IR data obtained for solid ${\rm IrF}_6$ /Xe (Tables 2 and 3) and the near IR-visible data for gas phase ${\rm IrF}_6$ /Xe allow a simpler, though perhaps more approximate, alternative method to be utilized.

The fact that there is no evidence of the IrF_6 -Xe complex in the room temperature vapor indicates that the 1:1 complex is not strongly bound. The observed frequency shifts, increased linewidths (see Table 3), and intensity enhancements in the 0.1% IrF_6 /Xe solid samples indicate that there is significant CT interaction between IrF_6 and Xe in the ground state. However, the solid state data pertain to an IrF_6 ·12Xe complex, whereas information on the 1:1 complex is of more intrinsic interest. Eq. (1) and the assumption that the complex is weak enough to allow n' to be set equal to 12 provides an approximate relationship between the 1:12 and the 1:1 complexes.

The stabilization energy of the ground state $[12X_0(r_8(^4A_2))]$ of the $IrF_6\cdot 12Xe$ complex can be estimated by assuming that Eq. (2) can be applied to the frequency shift data for the $r_8(^2T_1)$ state to find $\beta_0(r_8(^2T_1))$, and that $\beta_0(r_8(^4A_2))$ is roughly the same since both states stem from the $(t_{2g})^3$ configuration. These considerations lead to a value for $12X_0(r_8(^4A_2))$ of 300 cm⁻¹. The stabilization energy of the ground state $X_0(r_8(^4A_2))$ of the 1:1 complex is then 300 cm⁻¹. One would certainly expect this to be a lower limit since saturation effects have not been taken into account. The original assumption that the ground state complex is weak thus appears well justified.

The general picture that emerges for these weak charge transfer complexes

between Xe and MF $_6$ molecules is then as follows. The ground state is neutral probably with a shallow broad potential minimum somewhere near r° -4Å, the Xe-Xe approximate distance in a crystal lattice or a liquid. It is possible that there are a few vibrational quanta in this well but this is not a necessary condition imposed by our data. The excited state potential well is much deeper, more narrow, and the potential minimum is such that $r^*(\text{MF}_6^-\text{Xe}^+) < r^\circ(\text{MF}_6\text{Xe})$. These considerations also account nicely for the very broad (>10,000 cm $^{-1}$) Franck-Condon envelope observed for the CT transitions in all systems.

VI. CONCLUSION

The new electronic transitions which appear when certain transition metal hexafluorides are dissolved in liquid Xenon can be assigned as intermolecular charge-transfer transitions. The concomitant charge-transfer complexes are weakly bound.

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Table 1. Onset frequency for inter- and intramolecular charge-transfer transitions $[\sigma_{CT}(MF_6/Xe), \sigma_{CT}(MF_6)]$ and estimated electron affinities (E_a) for MF₆. A literature value for E_a (WF₆) is used in conjunction with Eq. (3) to determine the other E_a (MF₆) (see text).

MF ₆	σ _{CT} (MF ₆ /Xe) (cm ⁻¹)	σ _{CT} (MF ₆) (cm ⁻¹)	E _a (kcal/mole)
IrF ₆	9,800	20,000	181
$^{ReF}_{\boldsymbol{\epsilon}}$	>22,000	22,000	<146
WF ₆	36,600	52,000	(104) ^(a)
MoF ₆	26,300	50,000	133
UF ₆	>26,000	26,000	<134

⁽a) C. D. Cooper, R. N. Compton and P. W. Reinhardt, "Abstracts of Papers of the IXth International Conference on the Physics of Electronic and Atomic Collisons", edited by J. S. Risley and R. Geballe (University of Washington Press, Seattle, 1975), Vol. 2, p. 922.

Table 2. Observable near IR transitions of 0.1% $\rm IrF_6/Xe$. Frequencies are determined to $\pm 1~\rm cm^{-1}$.

MF ₆	Vacuum Wavenumber (cm ⁻¹)	FWHH (cm ⁻¹) ^(a)	Assignment
IrF ₆	5724	205	г ₈ (² т ₁)
	6016	168	$r_8 (^2T_1) + b.v.^{(b)}$
	6404	-	$r_8 (^2T_1) + s.v.^{(c)}$
	6641	-	r ₈ (² T ₁) + b.v.+ s.v.
	7814	203	r ₈ (² E)
	8036	•	r ₈ (² E) + b.v.
	8324	207	r ₆ (² T ₁)
	8575	-	$r_6 (^2 T_1) + b.v.$

⁽a) FWHH = full width at half height.

⁽b) b.v. = bending vibrations (v_4 , v_5 , v_6).

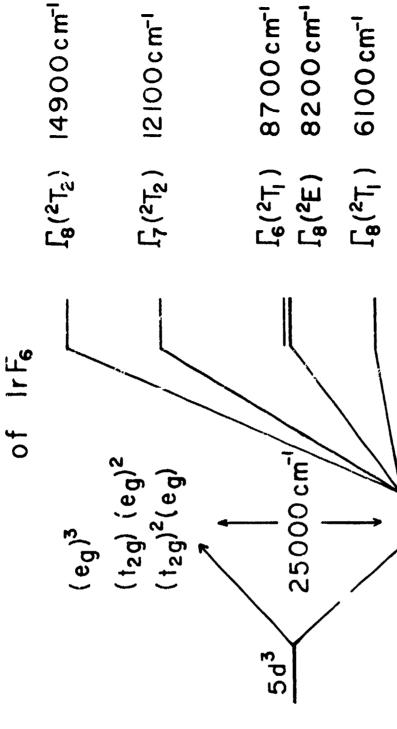
⁽c) s.v. = stretching vibrations (v_1, v_2, v_3) .

Table 3. Shifts in near IR origins between neat $\rm IrF_6$ and 0.1% $\rm IrF_6/Xe$.

	Neat IrF ₆ (cm ⁻¹)	0.1% IrF ₆ /Xe (cm ⁻¹)		Δ (cm ⁻¹)
$r_8 (^2 \tau_1)$	6114	5724		-390
r ₈ (² E)	8177	7814	,	-363
$r_6 (^2T_1)$	8701	8324		-377

Figure 1. Energy levels of IrF_6 . Rigorous symmetry labels (0^*_h) for each electronic state are given, along with the cubic Russell-Saunders state which correlates with the state for vanishing spin-orbit coupling. Since all the final states are gerade, the g label has been omitted in the right-hand column.

Energy Level Diagram



 $\Gamma_{\mathbf{g}}(^{4}A_{2}) \quad \text{Ocm}^{-1}$

(t₂g)³

 $(H^{20} + H^{2})$

+ H_{CF}

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